

Introduction

Hydrogels and their intrinsic instabilities have a wide variety of applications due to their high levels of hydrophilicity, biocompatibility, and stimulus-responsive behaviours against light, temperature, chemical, and electric fields, including:

- Actuators
- Stretchable electronics
- Tissue engineering
- Biosensing
- Drug delivery
- Tunable wettability or adhesiveness ...



The objective of this work includes three novel studies:

- develop a mixed & stable computational tool based on isogeometric analysis,
- provide an efficient framework for stability analysis of hydrogels,
- present numerical examples to show the robustness of the formulation.

Hydrogel Formulation

Coupled solid-deformation & fluid-diffusion

The governing equations are derived through the variational principle. Through the incompressibility constraint $det(\mathbf{F}) = 1 + \nu C$ and $p := \mu/\nu$, we consider the mixed incremental potential;

$$\Pi(\boldsymbol{\varphi}, p) = \int_{\mathcal{B}} \left[\psi - p(J - J_n) - \Delta t \, \phi^{\mathsf{dissipation}}(-\nabla p; \boldsymbol{F}_n) \right] \, dV$$

with the coupled deformation φ ($F = \nabla \varphi$) and fluid pressure p. C: concentration of solvent molecules, ν : volume of a solvent molecule. The stationarity of the potential implies

$$D\Pi[\delta\boldsymbol{\varphi}] = \int_{\mathcal{B}} \left[(\partial_{\boldsymbol{F}} \boldsymbol{\psi} - p \ J\boldsymbol{F}^{-T}) : \delta\boldsymbol{F} \right] dV - D\Pi^{\mathsf{EXT}}[\boldsymbol{\varphi}] = 0,$$

$$D\Pi[\delta p] = \int_{\mathcal{B}} \left[-(J - J_n)\delta p - \Delta t \partial_{\nabla p} \boldsymbol{\phi}^{\mathsf{dissipation}} \cdot \nabla \delta p \right] dV - D\Pi$$

Material model & solvent diffusion:

We adopt the free energy in [3]:

$$\psi(\mathbf{F}) = \underbrace{\frac{1}{2}Nk_BT(\mathbf{F}:\mathbf{F}-3-2\ln(J))}_{\text{stretching part}} \underbrace{-\frac{k_BT}{\nu}\left[(J-1)\ln(\frac{J}{J-1}) + \frac{\chi}{J}\right]}_{\text{mixing part}}$$

with N polymeric chains per volume, k_B Boltzmann factor, T temperature, χ polymer-solvent interaction parameter, D coefficient of diffusion. The Cauchy stress & the fluid flux:

$$oldsymbol{\sigma} = rac{1}{J} \mathbf{P} oldsymbol{F}^T$$
 and $\mathbf{q} = \partial_{
abla p} \phi^{\mathsf{dissipation}} = -rac{J_n - 1}{J_n} rac{J_n}{k_I}$

A Systematic Study of Swelling-Induced Instabilities of Hydrogels

Berkin Dortdivanlioglu & Christian Linder



Global and local instability modes²

PSfrag replacements

 $-\Pi^{\mathsf{EXT}}(\bar{\boldsymbol{\varphi}},\bar{p})$

 $\Pi^{\mathsf{EXT}}[\delta p] = 0 \ .$

 $\frac{D\nu}{BT}\nabla p.$

fluid pressure are interpolated independently using higher order NURBS.

$$\boldsymbol{\varphi}^{h} = \sum_{I=1}^{n_{en}^{\boldsymbol{\varphi}}} N_{I}^{\boldsymbol{\varphi}} \mathbf{u}_{I} \text{ and } p^{h} = \sum_{I=1}^{n_{en}^{p}} N_{I}^{pres} p_{I}$$





$$\underbrace{\begin{bmatrix} \boldsymbol{K}_{\boldsymbol{\varphi}\boldsymbol{\varphi}} & \boldsymbol{K}_{\boldsymbol{\varphi}p} \\ \boldsymbol{K}_{p\boldsymbol{\varphi}} & \boldsymbol{K}_{pp} \end{bmatrix}}_{\boldsymbol{K}} \begin{bmatrix} \Delta \boldsymbol{u} \\ \Delta p \end{bmatrix} =$$

$$oldsymbol{K}^* = oldsymbol{K}_{oldsymbol{arphi} arphi} - oldsymbol{K}_{pp} oldsymbol{K}_{pp}^{-1} oldsymbol{K}_{p oldsymbol{arphi}}$$
 and $oldsymbol{K}^{**} = -oldsymbol{H}$

Numerical Examples



Part II: Thin hydrogel ring immersed in water ($\lambda_0 = 1.1$)

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[7] B. Dortdivanlioglu & C Linder (2019). J. Mech. Phys. Solids 125:38-52. [8] B. Dortdivanlioglu et al. (2021). *J. Elast.* 145(1):31-47.

> Stanford University

Stanford DOERR School of Sustainability